Topology in Analysis

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The weak derivative of a locally intgegrable function $u : B \to \mathbb{R}^2$ is $\nabla u:B\to\mathbb{R}^{2\times 2}$ satisfying

$$
\int_B u \cdot \nabla \varphi \, dx = -\int_B \varphi \cdot \nabla u \, dx \qquad \forall \varphi \in C_c^{\infty}(B; \mathbb{R}^2)
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A function $u:B\to\mathbb{R}^2$ is in $W^{1,2}(B;\mathbb{R}^2)$ if it is weakly differntiable and $\left\| u \right\|_2$, $\left\| \nabla u \right\|_2$ are both finite.

Energy Functional & Admissibles

The Dirichlet energy functional $\mathbb{D}(u)$ can be used to describe the strain energy density under a deformation u .

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\mathbb{D}(u) = \int_B |\nabla u|_{\mathsf{F}}^2 \, \mathrm{d}x \qquad u \in W^{1,2}(B;\mathbb{R}^2)
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We shall consider the following space of admissible maps

$$
\mathcal{A} = \{u \in u_2 + W_0^{1,2}(B; \mathbb{R}^2) : \det \nabla u = 1 \text{ a.e. }\}
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This is commonly referred to as an **incompressibility** constraint. The chosen boundary condition is the **double covering map** $u_2 : B \to B' := \frac{1}{\sqrt{2}}$ $\frac{1}{2} B$ which is given (in complex coordinates) by

$$
r \mapsto \frac{1}{\sqrt{2}}r \qquad \theta \mapsto 2\theta
$$

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Plot of u_2

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\mathcal{V}(\Omega) := \{ \varphi \in \mathrm{id} + W_0^{1,2}(\Omega; \mathbb{R}^2) : \det \nabla \varphi = 1 \text{ a.e. } \}
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Then a mixed variation of u_2 is a map of the form

 $\psi \circ u_2 \circ \varphi \in \mathcal{A} \qquad \varphi \in \mathcal{V}(B) \quad \psi \in \mathcal{V}(B')$

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But just how big/small is this class?

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Consider a map $u \in C^1(B;\mathbb{R}^2)$ and a regular point $y \not\in u(\partial B)$. The **topological degree** is given by

$$
\deg(u, B, y) = \sum_{x \in u^{-1}\{y\}} \operatorname{sgn} \det \nabla u(x)
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We also have that the degree is constant on connected components of $u(B)$.

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However, we know more about u_2 .

- \bullet It is 1:1 at the origin.
- \bullet It is 2:1 everywhere else.

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What about a general admissible map?

$$
\left|u^{-1}\lbrace y\rbrace\right| \leq \deg(u, B, y) \quad \forall y
$$

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since admissibles preserve orientation.

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We then define the **associated self-map** $id_u : B \rightarrow B$ by

$$
id_u(x) = \begin{cases} x' & u^{-1}(u(x)) = \{x, x'\} \\ x & u^{-1}(u(x)) = \{x\} \end{cases}
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How many are there?

Lefschetz-Hopf Theorem

The Lefschetz number is

$$
\Lambda_{\mathrm{id}_u} = \sum_{k\geq 0} (-1)^k \operatorname{tr}((\mathrm{id}_u)_* | H_k(B; \mathbb{Q}))
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The Lefschetz-Hopf theorem for ENR's states that

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\Lambda_{\mathsf{id}_u} = \sum_{x \in \mathsf{Fix}(\mathsf{id}_u)} i(\mathsf{id}_u, x) = \sum_{x \in \mathsf{Fix}(\mathsf{id}_u)} 1 = |\mathsf{Fix}(\mathsf{id}_u)|
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Result: Any orientation preserving injective map self-map of the ball has either 1 fixed point or infinite fixed points.

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We now know that any $u \in \mathcal{A}$ such that $\left|u^{-1}\{y\}\right| < \infty$ satisfies

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- \bullet It is 1:1 somewhere in its domain.
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- \bullet It is 2:1 everywhere else.

This is interesting when compared to the C^1 case.

If $u \in C^1(X;Y)$ is proper for X connected and Y simply connected, then

$$
\det \nabla u \neq 0 \Rightarrow u \text{ is a diffeomorphism}
$$

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so a discontinuity in the derivative is essential for 2:1 behaviour.

What if there exists y such that $\left|u^{-1}\{y\}\right| = \infty$? Either countable or uncountable?

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In the uncountable case, we could have a line in B that is compressed to a point in the image.

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We could stitch together a sequence of maps that contract curves by progressively larger amounts.

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We could stitch together a sequence of maps that contract curves by progressively larger amounts.

We consider maps on an annulus $B_{r_n} - B_{r_{n+1}}$ of the form

$$
u_n(x) = \ell_n(x)R(\omega_n(x))x
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Here ℓ_n are scalar functions governing the contracting factor and ω_n are scalar functions governing rotation.

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The r_n must be a decreasing sequence in [0, 1] that is bounded away from zero and has initial condition $r_0 = 1$.

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We calculate the Jacobian using polar coordinates

$$
\det \nabla u_n = (r\ell_{n,r} + \ell_n)\ell_n(\omega_{n,\theta} + 1) - r\ell_{n,\theta}\ell_n^2\omega_{n,r}
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Derived PDE

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If we assume $\ell_n = \ell_n(r)$ then the Jacobian constraint is just

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$$
\Rightarrow \omega = \left(\frac{1}{(r\ell_{n,r} + \ell_n)\ell_n} - 1\right)\theta + C_n(r)
$$

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for arbitrary functions C_n , which we shall set to zero.

For u to be continuous, we must have that a change in a multiple of 2π in θ results in a change of a multiple of 2π in ω .

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This leads us to the following ODE for ℓ_n :

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(r\ell_{n,r}+\ell_n)\ell_n=1
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(r\ell_{n,r}+\ell_n)\ell_n=1
$$

which has solution

$$
\ell_n(r) = \frac{\sqrt{c_n + r^2}}{r}
$$

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for some constants c_n .

Continuity & Contracting Factors

To ensure continuity we must have

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\ell_{n+1}(r_n)=\ell_n(r_n)
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and to get increasing contractions we must have an increasing sequence α_n with $\alpha_1 = 1$ such that

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\ell_n(r_n)=\frac{r_n}{\alpha_n}
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$$

Combining these with the equation for ℓ_n in the previous slide, we obtain

$$
\ell_n(r) = \frac{1}{r} \sqrt{r_n^2 \left(\left(\frac{r_n}{\alpha_n} \right)^2 - 1 \right) + r^2}
$$

$$
r_n = \frac{\alpha_n}{\sqrt{2}} \sqrt{1 \pm \sqrt{1 + \left(\frac{2}{\alpha_n \alpha_{n-1}} \right)^2 (\alpha_{n-1}^2 - r_{n-1}^2) r_{n-1}^2}, \quad r_0 = 1
$$

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How can we modify this approach to successfully construct a solution?

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Change from annulus to sector based partitioning?

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Make the ℓ_n depend on θ ?

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How can we modify this approach to successfully construct a solution?

Change from annulus to sector based partitioning?

Make the ℓ_n depend on θ ?

Try to contract a ray rather than a circle?